## AME40541/60541: Finite Element Methods Homework 5: Due Monday, March 8, 2021

**Problem 1:** (35 points) (AME 60541 only) Consider a system of m second-order conservation laws in a d-dimensional domain  $\Omega \subset \mathbb{R}^d$ 

$$\nabla \cdot F(U, \nabla U) = S(U, \nabla U) \quad \text{in } \Omega,$$

where  $U: \Omega \to \mathbb{R}^m$  is the state,  $F(U, \nabla U) \in \mathbb{R}^{m \times d}$  is the flux function (operator), and  $S(U, \nabla U) \in \mathbb{R}^m$  is a source term. The boundary conditions are  $U = \overline{U}$  on  $\partial \Omega_D$  and  $F(U, \nabla U)n = \overline{q}$  on  $\partial \Omega_N$ , where  $\overline{U}: \partial \Omega_D \to \mathbb{R}^m$  and  $\overline{q}: \partial \Omega_N \to \mathbb{R}^m$  are known boundary functions,  $\partial \Omega = \overline{\partial \Omega_D \cup \partial \Omega_N}$ , and  $n \in \mathbb{R}^d$  is the outward normal.

- (a) (5 points) Write the conservation law in indicial notation. Drop the arguments to the flux function and source term.
- (b) (10 points) Construct both the weighted residual and weak formulation of the governing equations.
- (c) (5 points) What conditions must an approximate solution  $U_h(x) \approx U(x)$  satisfy if applying (1) the method of weighted residuals or (2) the Ritz method? Why is it difficult to construct  $U_h$  if F is nonlinear in U or  $\nabla U$  if using the method of weighted residuals?
- (d) (15 points, 5 each) Write each of the following PDEs as a general conservation law, i.e., identify the state (U), flux function  $(F(U, \nabla U))$ , source term  $(S(U, \nabla U))$ , and boundary conditions  $(\overline{U}, \overline{q})$ . Also identify the PDE as linear or nonlinear (justify your answer) and the number of solution components (m).
  - The second-order, linear elliptic PDE over the domain  $\Omega \subset \mathbb{R}^d$

$$(k_{ij}u_{,j})_{,i} = f \text{ in } \Omega, \qquad u = \bar{u} \text{ on } \partial\Omega_D, \qquad k_{ij}u_{,j}n_i = \bar{t} \text{ on } \partial\Omega_N,$$

where  $u: \Omega \to \mathbb{R}$  is the unknown solution,  $k: \Omega \to M_{n,n}(\mathbb{R})$  are the elliptic coefficients,  $f: \Omega \to \mathbb{R}$ is the source term,  $\bar{u}: \partial\Omega_D \to \mathbb{R}$  and  $\bar{t}: \partial\Omega_N \to \mathbb{R}$  are boundary conditions,  $n: \partial\Omega \to \mathbb{R}^d$  is the outward normal, and the boundary is  $\partial\Omega = \overline{\partial\Omega_D \cup \partial\Omega_N}$ .

• The linear elasticity equations govern the deformation of a domain  $\Omega \subset \mathbb{R}^d$  subject to loads  $f: \Omega \to \mathbb{R}^d$ 

 $\sigma_{ij,j} + f_i = 0$  in  $\Omega$ ,  $u_i = \bar{u}_i$  on  $\partial \Omega_D$ ,  $\sigma_{ij} n_j = \bar{t}_i$  on  $\partial \Omega_N$ 

where  $\sigma: \Omega \to M_{d,d}(\mathbb{R})$  is the stress field,  $\bar{u}: \partial\Omega_D \to \mathbb{R}^d$  is the prescribed boundary displacement field and  $\bar{t}: \partial\Omega_N \to \mathbb{R}^d$  is the prescribed traction,  $n: \partial\Omega \to \mathbb{R}^d$  is the outward normal, and the boundary is  $\partial\Omega = \overline{\partial\Omega_D \cup \partial\Omega_N}$ . The stress is related to the strain field  $\epsilon: \Omega \to M_{d,d}(\mathbb{R})$  using Hooke's law (linear elastic material) and the strains are assumed infinitesimal

$$\sigma_{ij} = C_{ijkl}\epsilon_{kl}, \qquad \epsilon_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i}).$$

• The incompressible Navier-Stokes equations (see equations in Problem 2)

**Problem 2:** (20 points) Consider the incompressible Navier-Stokes equations that govern the flow of an incompressible fluid with density  $\rho : \Omega \to \mathbb{R}_{>0}$  and viscosity  $\nu : \Omega \to \mathbb{R}_{>0}$  through a domain  $\Omega \subset \mathbb{R}^3$ 

$$\begin{aligned} -\frac{\partial}{\partial x_1} \left(\rho \nu \frac{\partial v_1}{\partial x_1}\right) - \frac{\partial}{\partial x_2} \left(\rho \nu \frac{\partial v_1}{\partial x_2}\right) - \frac{\partial}{\partial x_3} \left(\rho \nu \frac{\partial v_1}{\partial x_3}\right) + \rho v_1 \frac{\partial v_1}{\partial x_1} + \rho v_2 \frac{\partial v_1}{\partial x_2} + \rho v_3 \frac{\partial v_1}{\partial x_3} + \frac{\partial p}{\partial x_1} = 0 \\ -\frac{\partial}{\partial x_1} \left(\rho \nu \frac{\partial v_2}{\partial x_1}\right) - \frac{\partial}{\partial x_2} \left(\rho \nu \frac{\partial v_2}{\partial x_2}\right) - \frac{\partial}{\partial x_3} \left(\rho \nu \frac{\partial v_2}{\partial x_3}\right) + \rho v_1 \frac{\partial v_2}{\partial x_1} + \rho v_2 \frac{\partial v_2}{\partial x_2} + \rho v_3 \frac{\partial v_2}{\partial x_3} + \frac{\partial p}{\partial x_2} = 0 \\ -\frac{\partial}{\partial x_1} \left(\rho \nu \frac{\partial v_3}{\partial x_1}\right) - \frac{\partial}{\partial x_2} \left(\rho \nu \frac{\partial v_3}{\partial x_2}\right) - \frac{\partial}{\partial x_3} \left(\rho \nu \frac{\partial v_3}{\partial x_3}\right) + \rho v_1 \frac{\partial v_3}{\partial x_1} + \rho v_2 \frac{\partial v_3}{\partial x_2} + \rho v_3 \frac{\partial v_3}{\partial x_3} + \frac{\partial p}{\partial x_3} = 0 \\ \frac{\partial v_1}{\partial x_1} + \frac{\partial v_2}{\partial x_2} + \frac{\partial v_3}{\partial x_3} = 0 \end{aligned}$$

where  $v : \Omega \to \mathbb{R}^3$  with components  $v = (v_1, v_2, v_3)$  is the velocity of the fluid and  $p : \Omega \to \mathbb{R}_{>0}$  is the pressure. The boundary is partitioned into two pieces:  $\partial \Omega = \overline{\partial \Omega_D \cup \partial \Omega_N}$ , where  $n : \partial \Omega \to \mathbb{R}^3$  is the outward unit normal. The flow velocity and pressure are prescribed along  $\partial \Omega_D$ 

$$v = \bar{v}, \ p = \bar{p} \quad \text{on } \partial\Omega_D,$$

where  $\bar{u}: \partial\Omega_D \to \mathbb{R}^3$  is the prescribed velocity with components  $\bar{v} = (\bar{v}_1, \bar{v}_2, \bar{v}_3)$  and  $\bar{p}: \partial\Omega_D \to \mathbb{R}_{>0}$  is the prescribed pressure. The traction is prescribed as  $\bar{t} = (\bar{t}_1, \bar{t}_2, \bar{t}_3)$  along  $\partial\Omega_N$ 

$$\rho\nu\left(\frac{\partial v_1}{\partial x_1}n_1 + \frac{\partial v_1}{\partial x_2}n_2 + \frac{\partial v_1}{\partial x_3}n_3\right) - pn_1 = \rho\bar{t}_1$$
  

$$\rho\nu\left(\frac{\partial v_2}{\partial x_1}n_1 + \frac{\partial v_2}{\partial x_2}n_2 + \frac{\partial v_2}{\partial x_3}n_3\right) - pn_2 = \rho\bar{t}_2 \quad \text{on } \partial\Omega_N$$
  

$$\rho\nu\left(\frac{\partial v_3}{\partial x_1}n_1 + \frac{\partial v_3}{\partial x_2}n_2 + \frac{\partial v_3}{\partial x_3}n_3\right) - pn_3 = \rho\bar{t}_3.$$

Re-write these equations in indicial notation and construct the weak form of the equations. Observe that in this form the equations can easily be generalized from three dimensions to d dimensions. This means you have also derived the weak formulation of the 1d, 2d (and higher dimensions!) incompressible Navier-Stokes equations as well.

**Problem 3:** (40 points) Consider the equations associated with a simply supported beam and subjected to a uniform transverse load  $q = q_0$ :

$$\frac{d^2}{dx^2} \left( EI \frac{d^2 w}{dx^2} \right) = q_0 \quad \text{for} \quad 0 < x < L$$
$$w = EI \frac{d^2 w}{dx^2} = 0 \quad \text{at} \quad x = 0, L.$$

Take L = 1, EI = 1, and  $q_0 = -1$  and approximate the solution using the following methods:

- (a) the method of weighted residuals (Galerkin) using the two-term trigonometric basis  $w(x) \approx w_2(x) = c_1 \sin\left(\frac{\pi x}{L}\right) + c_2 \sin\left(\frac{2\pi x}{L}\right)$ ,
- (b) the method of weighted residuals (collocation) using the same trigonometric basis and collocation nodes  $x_1 = 0.25$  and  $x_2 = 0.75$ ,
- (c) the Ritz method using the same trigonometric basis, and
- (d) the Ritz method using a two-term polynomial basis  $(w(x) \approx w_2(x) = c_1 x(x-L) + c_2 x^2 (x-L)^2)$ .

For each method, verify the solution basis satisfies the appropriate conditions. Plot the approximate solution generated by each method as well as the analytical solution. In a separate figure, plot the error of each method  $e(x) = |w(x) - \tilde{w}(x)|$ , where  $\tilde{w}$  is the approximate solution, and the residual over the domain. Finally, quantify the error of each approximation using the  $L^2$ -norm

$$e_{L^2(\Omega)} = \int_{\Omega} |e(x)|^2 \, dV.$$

I recommend using some symbolic mathematics software (Maple, Mathematica, MATLAB, etc) to assist with the calculations.